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# Exact Relationship Between Compressibility of Solution and Compressibilities of its Pure Constituents

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It has been shown that  $\bar{V}_1(T, p, x_1) = \bar{V}_1^0(T, p - \pi_1)$  and that  $\bar{V}_2(T, p, x_2) = \bar{V}_2^0(T, p - \pi_2)$ , where  $\pi_1$  is the osmotic pressure of the solvent in a binary solution and  $\pi_2$  is the osmotic pressure of the solution solute. An increase in pressure applied to the solution, the pure solvent and pure solute from  $p$  to  $p + \Delta p$  compresses the solution solvent and pure solvent by equal fractions

$$-\frac{\Delta \bar{V}_1(T, p, x_1)}{\bar{V}_1(T, p, x_1)} = -\frac{\Delta \bar{V}_1^0(T, p - \pi_1)}{\bar{V}_1^0(T, p - \pi_1)}$$

and compresses the solution solute and pure solute by equal fractions

$$-\frac{\Delta \bar{V}_2(T, p, x_2)}{\bar{V}_2(T, p, x_2)} = -\frac{\Delta \bar{V}_2^0(T, p - \pi_2)}{\bar{V}_2^0(T, p - \pi_2)}$$

Thus, the compressibilities of the solution solvent and pure solvent are related as

$$\kappa_1(T, p, x_1) = \kappa_1^0(T, p - \pi_1);$$

and likewise, the solution solute and pure solute compressibilities are related as

$$\kappa_2(T, p, x_2) = \kappa_2^0(T, p - \pi_2).$$

Since

$$-\frac{\Delta V}{V} = -n_1 \frac{\Delta \bar{V}_1}{\bar{V}_1} - n_2 \frac{\Delta \bar{V}_2}{\bar{V}_2},$$

it is shown that the compressibility of the solution and the compressibilities of the pure constituents are related as

$$\kappa_{12} V = n_1 \bar{V}_1^0 \kappa_1^0 + n_2 \bar{V}_2^0 \kappa_2^0$$

where  $\kappa_{12}$  and  $V$  are measured at  $p$ ,  $\kappa_1^0$  and  $\bar{V}_1^0$  are measured at  $p - \pi_1$ ,  $\kappa_2^0$  and  $\bar{V}_2^0$  are measured at  $p - \pi_2$  and all are measured at the same  $T$ .

## INTRODUCTION

The purpose of this article is to derive an exact relationship between the compressibility of a binary solution,  $\kappa_{12}$ , and the compressibilities of its pure constituents,  $\kappa_1^0$  and  $\kappa_2^0$ . The derivation will be based on the recognition that all species in a solution enhance the tension in the coupling forces, and thereby increase the partial molar volume of each species in the solution.<sup>1,2</sup>

### Compressibilities of solution, solvent, and solute

A binary solution contains  $n_1$  moles of solvent and  $n_2$  moles of solute. The volume of the solution at  $T$  and applied pressure  $p$  is

$$V(T, p, n_1, n_2) = n_1 \bar{V}_1(T, p, n_1, n_2) + n_2 \bar{V}_2(T, p, n_1, n_2), \quad (1)$$

where  $\bar{V}_1(T, p, n_1, n_2)$  is the partial molar volume of the solvent (i.e.  $\partial V/\partial n_1|T, p, n_2$ ) and  $\bar{V}_2(T, p, n_1, n_2)$  is partial molar volume of the solute (i.e.  $\partial V/\partial n_2|T, p, n_1$ ). If  $T$  is held constant as the applied pressure increases from  $p$  to  $p + \Delta p$ , the solution volume decreases to

$$V(T, p + \Delta p, n_1, n_2) = n_1 \bar{V}_1(T, p + \Delta p, n_1, n_2) + n_2 \bar{V}_2(T, p + \Delta p, n_1, n_2). \quad (2)$$

The decrease in solution volume  $-\Delta V$  due to the increase in pressure  $\Delta p$  is

$$-\Delta V(T, p, n_1, n_2) = V(T, p + \Delta p, n_1, n_2) - V(T, p, n_1, n_2). \quad (3)$$

The partial molar volumes of the solvent and solute also decrease by the amounts

$$-\Delta \bar{V}_1(T, p, n_1, n_2) = \bar{V}_1(T, p + \Delta p, n_1, n_2) - \bar{V}_1(T, p, n_1, n_2) \quad (4)$$

and

$$-\Delta \bar{V}_2(T, p, n_1, n_2) = \bar{V}_2(T, p + \Delta p, n_1, n_2) - \bar{V}_2(T, p, n_1, n_2). \quad (5)$$

By subtracting Eq. (1) from Eq. (2) and using Eqs. (3), (4), and (5), it follows that

$$-\Delta V(T, p, n_1, n_2) = -n_1 \Delta \bar{V}_1(T, p, n_1, n_2) - n_2 \Delta \bar{V}_2(T, p, n_1, n_2). \quad (6)$$

Dividing Eq. (6) by  $V$ ,

$$-\frac{\Delta V}{V} = -n_1 \frac{\Delta \bar{V}_1}{V} - n_2 \frac{\Delta \bar{V}_2}{V}. \quad (7)$$

The compressibility of the solution is defined as

$$\kappa_{12}(T, p, n_1, n_2) = -\frac{\Delta V}{V} \Delta p^{-1} \quad (8)$$

and is a function of  $T, p, n_1$  and  $n_2$ . Likewise, the compressibility of the solution solvent can be defined as

$$\kappa_1(T, p, n_1, n_2) = -\frac{\Delta\bar{V}_1}{\bar{V}_1} \Delta p^{-1}, \quad (9)$$

and the compressibility of the solution solute can be defined as

$$\kappa_2(T, p, n_1, n_2) = -\frac{\Delta\bar{V}_2}{\bar{V}_2} \Delta p^{-1}. \quad (10)$$

Combining Eqs. (8), (9) and (10) with Eq. (7),

$$\begin{aligned} \kappa_{12}(T, p, n_1, n_2) &= \frac{n_1 \bar{V}_1(T, p, n_1, n_2)}{V(T, p, n_1, n_2)} \kappa_1(T, p, n_1, n_2) \\ &+ \frac{n_2 \bar{V}_2(T, p, n_1, n_2)}{V(T, p, n_1, n_2)} \kappa_2(T, p, n_1, n_2). \end{aligned} \quad (11)$$

The solution was constituted by combining  $n_1^0$  moles of pure solvent, having a volume  $V_1^0(T, p) = n_1^0 \bar{V}_1^0(T, p)$ , with  $n_2^0$  moles of pure solute, having a volume  $V_2^0(T, p) = n_2^0 \bar{V}_2^0(T, p)$ .  $\bar{V}_1^0$  and  $\bar{V}_2^0$  are the molar volumes of the pure solvent and pure solute, respectively, and both are functions of  $T$  and  $p$ . Increasing the pressure applied to the pure solvent from  $p$  to  $p + \Delta p$  at constant  $T$  decreases its molar volume an amount

$$-\Delta\bar{V}_1^0(T, p) = \bar{V}_1^0(T, p + \Delta p) - \bar{V}_1^0(T, p)$$

and its compressibility is

$$\kappa_1^0(T, p) = -\frac{\Delta\bar{V}_1^0}{\bar{V}_1^0} \Delta p^{-1}. \quad (12)$$

Likewise, increasing the pressure applied to the pure solute from  $p$  to  $p + \Delta p$  decreases its molar volume an amount

$$-\Delta\bar{V}_2^0(T, p) = \bar{V}_2^0(T, p + \Delta p) - \bar{V}_2^0(T, p)$$

and the compressibility of pure solute is

$$\kappa_2^0(T, p) = -\frac{\Delta\bar{V}_2^0}{\bar{V}_2^0} \Delta p^{-1}. \quad (13)$$

### Relationship between compressibilities of solution solvent $\kappa_1$ and pure solvent $\kappa_1^0$

The chemical potential of either species  $i$  in a binary solution can be altered by changing  $T, p$  or the molar amount of species  $i$  according to

$$d\mu_i(T, p, n_1, n_2) = -\bar{S}_i dT + \bar{V}_i dp + \frac{\partial\mu_i}{\partial x_i} dx_i,$$

where  $\bar{S}_i$  is the partial molar entropy of species  $i$  and  $x_i = n_i/(n_1 + n_2)$ , the mole fraction of species  $i$ . Increasing the pressure applied to the solution from an initial  $p_i$  to  $p_i + \Delta p$  at constant  $T$  and  $x_1$  increases the chemical potential of the solvent in the solution an amount

$$\mu_1(T, p_i + \Delta p, x_1) - \mu_1(T, p_i, x_1) = \int_{p_i}^{p_i + \Delta p} \bar{V}_1 dp. \quad (14)$$

For pure solvent,  $x_1^0 = 1$ . Decreasing the pressure applied to it from  $p_i$  to  $p_i - \Delta p$  at constant  $T$  decreases its chemical potential an amount

$$\mu_1^0(T, p_i - \Delta p) - \mu_1^0(T, p_i) = \int_{p_i}^{p_i - \Delta p} \bar{V}_1^0 dp. \quad (15)$$

$\bar{V}_1$  and  $\bar{V}_1^0$  are both functions of  $T$  and  $p$  but they are not equal at the same  $T$  and  $p$ . However, there is an exact relationship between them.

If  $\Delta p$  in Eqs. (14) and (15) is equal to the osmotic pressure of the solvent in the solution,  $\pi_1$ , then the sum of the left sides of Eqs. (14) and (15) equals zero since

$$\mu_1(T, p_i + \pi_1, x_1) = \mu_1^0(T, p_i) \quad (16)$$

and

$$\mu_1(T, p_i, x_1) = \mu_1^0(T, p_i - \pi_1). \quad (17)$$

That is, if the pressure applied to the solution and pure solvent differ by  $\pi_1$ , then the chemical potentials of the solution solvent and pure solvent must be equal at constant  $T$ . Thus, it follows that

$$\int_{p_i}^{p_i + \pi_1} \bar{V}_1 dp = \int_{p_i - \pi_1}^{p_i} \bar{V}_1^0 dp. \quad (18)$$

$\bar{V}_1$  and  $\bar{V}_1^0$  are state functions of  $T$  and  $p$  and vary continuously with  $p$ ; so  $p_i$  in Eq. (18) can have any value and it can be increased to  $p_i + \pi_1$ ; so again

$$\int_{p_i + \pi_1}^{p_i + 2\pi_1} \bar{V}_1 dp = \int_{p_i}^{p_i + \pi_1} \bar{V}_1^0 dp. \quad (19)$$

The sum of Eqs. (18) and (19) becomes

$$\int_{p_i}^{p_i + 2\pi_1} \bar{V}_1 dp = \int_{p_i - \pi_1}^{p_i + \pi_1} \bar{V}_1^0 dp. \quad (20)$$

Thus,  $\bar{V}_1$  and  $\bar{V}_1^0$  vary with  $p$  and they must vary so as to satisfy Eqs. (18), (19) and (20). Let us assume that  $\bar{V}_1(T, p + \pi_1, x_1) = \bar{V}_1^0(T, p)$  and show that Eqs. (18), (19) and (20) are satisfied only if, for every  $p$ ,  $\bar{V}_1(T, p + \pi_1, x_1)$  does equal  $\bar{V}_1^0(T, p)$ .

The variation of the molar volume of the pure solvent from  $p_i - \pi_1$  to  $p_i + \pi_1$  at constant  $T$  is illustrated as the curve labelled  $\bar{V}_1^0(T, p)$  in Figure 1. If our assumption is valid, then the curve labelled  $\bar{V}_1(T, p, x_1)$  represents the variation of the partial molar volume of the solution solvent from  $p_i$  to  $p_i + 2\pi_1$ . By comparing the curve for  $\bar{V}_1^0(T, p)$  from  $p_i - \pi_1$  to  $p_i + \pi_1$  with the curve for  $\bar{V}_1(T, p, x_1)$  from  $p_i$  to  $p_i + 2\pi_1$ , several facts are clear:

- 1)  $\bar{V}_1^0(T, p_i - \pi_1) = \bar{V}_1(T, p_i, x_1)$ ;
- 2)  $\bar{V}_1^0(T, p_i) = \bar{V}_1(T, p_i + \pi_1, x_1)$ ;
- 3)  $\bar{V}_1^0(T, p_i + \pi_1) = \bar{V}_1(T, p_i + 2\pi_1, x_1)$ ;
- 4) 
$$\int_{p_i - \pi_1}^{p_i} \bar{V}_1^0 dp = \int_{p_i}^{p_i + \pi_1} \bar{V}_1 dp$$
;

i.e. the area under  $\bar{V}_1^0(T, p)$  from  $p_i - \pi_1$  to  $p_i$  equals the area under  $\bar{V}_1(T, p, x_1)$  from  $p_i$  to  $p_i + \pi_1$  and Eq. (18) is satisfied;

5) Eqs. (19) and (20) are also satisfied.

Next, assume that the curve for the partial molar volume of the solvent is, as before,  $\bar{V}_1(T, p, x_1)$  in Figure 1, and assume that the molar volume of the pure solvent varies as the broken curve labelled  $\bar{V}_1^0(T, p)$  from  $p_i - \pi_1$  to  $p_i + \pi_1$ . These curves assume that the pure solvent is more compressible than

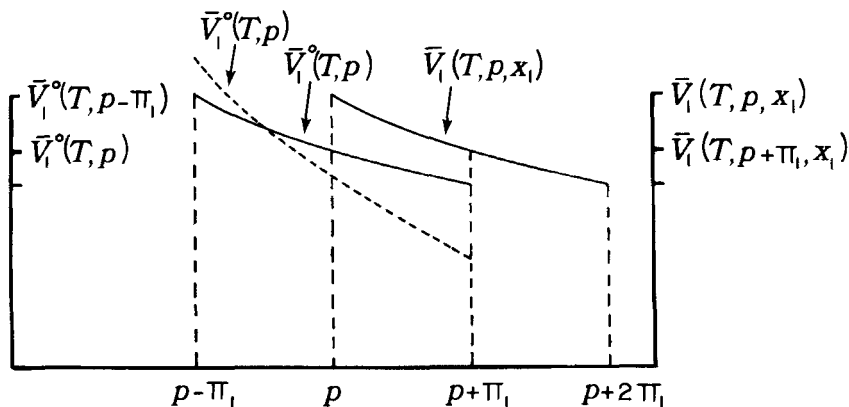


FIGURE 1 Molar volumes (ordinate) versus applied pressures (abscissa) of pure solvent and solution solvent. Solid curves demonstrate that the compressibility of pure solvent from  $p - \pi_1$  to  $p + \pi_1$  equals compressibility of solution solvent from  $p$  to  $p + 2\pi_1$  at the same  $T$ . Broken curve incorrectly assumes that compressibility of pure solvent differs from compressibility of solution solvent.

the solution solvent. The broken curve  $\bar{V}_1^0(T, p)$  was drawn to satisfy Eq. (18) but it can not also satisfy Eqs. (19) and (20), i.e.

$$\int_{p_i}^{p_i + \pi_1} \bar{V}_1^0 dp \neq \int_{p_i + \pi_1}^{p_i + 2\pi_1} \bar{V}_1 dp$$

and

$$\int_{p_i - \pi_1}^{p_i + \pi_1} \bar{V}_1^0 dp \neq \int_{p_i}^{p_i + 2\pi_1} \bar{V}_1 dp.$$

Since the compressibility of the pure solvent and  $\bar{V}_1^0(T, p)$  from  $p_i - \pi_1$  to  $p_i + \pi_1$  must satisfy Eqs. (18), (19) and (20), we can conclude that our first assumption was valid, i.e.  $\bar{V}_1^0(T, p) = \bar{V}_1(T, p + \pi_1, x_1)$ , and the compressibility of pure solvent at every  $p$  must be identical with the compressibility of the solution solvent at every  $p + \pi_1$  at constant  $T$  and  $x_1$ . It also follows that

$$\bar{V}_1(T, p, x_1) = \bar{V}_1^0(T, p - \pi_1) \quad (21)$$

and that

$$\Delta \bar{V}_1(T, p, x_1) = \Delta \bar{V}_1^0(T, p - \pi_1)$$

if the pressure applied to the pure solvent changes from  $p - \pi_1$  to  $p + \Delta p - \pi_1$  while the pressure applied to the solution changes from  $p$  to  $p + \Delta p$ . Thus,

$$\frac{\Delta \bar{V}_1(T, p, x_1)}{\bar{V}_1(T, p, x_1)} = \frac{\Delta \bar{V}_1^0(T, p - \pi_1)}{\bar{V}_1^0(T, p - \pi_1)}.$$

That is, the compressibility of the solution solvent at  $T$  and  $p$  must equal the compressibility of the pure solvent at  $T$  and  $p - \pi_1$

$$\kappa_1(T, p, x_1) = \kappa_1^0(T, p - \pi_1). \quad (22)$$

#### Relationship between compressibilities of solution solute $\kappa_2$ and pure solute $\kappa_2^0$

The same argument as above can be exploited to derive the relationship between the compressibilities of the solution solute and pure solute. Increasing the pressure applied to the solution from an initial pressure  $p_i$  to  $p_i + \Delta p$  increases the chemical potential of the solution solute an amount

$$\mu_2(T, p_i + \Delta p, x_2) - \mu_2(T, p_i, x_2) = \int_{p_i}^{p_i + \Delta p} \bar{V}_2 dp, \quad (23)$$

at constant  $T$  and  $x_2$ . For pure solute  $x_2^0 = 1$ . Decreasing the pressure applied to it from  $p_i$  to  $p_i - \Delta p$  at constant  $T$  decreases its chemical potential an amount

$$\mu_2^0(T, p_i - \Delta p) - \mu_2^0(T, p_i) = \int_{p_i}^{p_i - \Delta p} \bar{V}_2^0 dp. \tag{24}$$

$\bar{V}_2$  and  $\bar{V}_2^0$  are both functions of  $T$  and  $p$  and there is an exact relationship between them.

If  $\Delta p$  in Eqs. (23) and (24) is equal to the osmotic pressure of the solute in the solution,  $\pi_2$ , then the sum of left sides of Eqs. (23) and (24) equals zero since

$$\mu_2(T, p_i + \pi_2, x_2) = \mu_2^0(T, p_i) \tag{25}$$

and

$$\mu_2(T, p_i, x_2) = \mu_2^0(T, p_i - \pi_2). \tag{26}$$

Thus, it follows that

$$\int_{p_i}^{p_i + \pi_2} \bar{V}_2 dp = \int_{p_i - \pi_2}^{p_i} \bar{V}_2^0 dp. \tag{27}$$

It also follows that

$$\int_{p_i}^{p_i + 2\pi_2} \bar{V}_2 dp = \int_{p_i - \pi_2}^{p_i + \pi_2} \bar{V}_2^0 dp, \tag{28}$$

since  $\bar{V}_2$  and  $\bar{V}_2^0$  are state functions of  $T$  and  $p$ . Only if

$$\bar{V}_2(T, p, x_2) = \bar{V}_2^0(T, p - \pi_2) \tag{29}$$

are both Eqs. (27) and (28) satisfied. It also follows that

$$-\Delta \bar{V}_2(T, p, x_2) = -\Delta \bar{V}_2^0(T, p - \pi_2), \tag{30}$$

if the pressure applied to the solution increases from  $p$  to  $p + \Delta p$  while the pressure applied to the pure solute increases from  $p - \pi_2$  to  $p + \Delta p - \pi_2$ . Thus,

$$\frac{\Delta \bar{V}_2(T, p, x_2)}{\bar{V}_2(T, p, x_2)} = \frac{\Delta \bar{V}_2^0(T, p - \pi_2)}{\bar{V}_2^0(T, p - \pi_2)} \tag{31}$$

and

$$\kappa_2(T, p, x_2) = \kappa_2^0(T, p - \pi_2). \tag{32}$$

That is, the compressibility of the solution solute at  $T$  and  $p$  must equal the compressibility of the pure solute at  $T$  and  $p - \pi_2$ .



**Relationship between the compressibility of the solution  $\kappa_{12}$  and the compressibilities of the pure solvent  $\kappa_1^0$  and pure solute  $\kappa_2^0$**

The relationship between the compressibility of a binary solution  $\kappa_{12}(T, p, n_1, n_2)$  and the compressibilities of its solution solvent and solution solute was derived as Eq. (11). Eqs. (21), (22), (29) and (32) can be used to substitute for  $\bar{V}_1, \kappa_1, \bar{V}_2$  and  $\kappa_2$  in Eq. (11). Thus, Eq. (11) becomes

$$\kappa_{12}(T, p, n_1, n_2)V(T, p, n_1, n_2) = n_1\bar{V}_1^0(T, p - \pi_1)\kappa_1^0(T, p - \pi_1) + n_2\bar{V}_2^0(T, p - \pi_2)\kappa_2^0(T, p - \pi_2). \quad (33)$$

The product of the volume and compressibility of the solution at  $T, p, n_1$  and  $n_2$  must equal the product of the number of moles of solution solvent at  $T, p, n_1$  and  $n_2$ , the molar volume and the compressibility of pure solvent at  $T$  and  $p - \pi_1$  plus the product of the number of moles of solution solute, the molar volume and the compressibility of the pure solute at  $T$  and  $p - \pi_2$ .

## CONCLUSION

Equation (33) is an exact relationship between the compressibility and volume of a binary solution and the compressibilities and molar volumes of its pure constituents.  $\kappa_{12}$  and  $V_{12}$  can be measured at  $T$  and  $p$ ;  $\bar{V}_1^0$  and  $\kappa_1^0$  can be measured at  $T$  and  $p - \pi_2$ ;  $\bar{V}_2^0$  and  $\kappa_2^0$  can be measured at  $T$  and  $p - \pi_2$ . If  $n_1$  and  $n_2$  are known at  $T$  and  $p$ , then these data must satisfy Eq. (33). If Eq. (33) is not satisfied, then  $n_1$  and  $n_2$  were not known at  $T$  and  $p$ .

Eq. (33) can also be written

$$\begin{aligned} \kappa_{12}(T, p + \pi_1, n_1, n_2)V_{12}(T, p + \pi_1, n_1, n_2) \\ = n_1\bar{V}_1^0(T, p)\kappa_1^0(T, p) + n_2\bar{V}_2^0(T, p + \pi_1 - \pi_2)\kappa_2^0(T, p + \pi_1 - \pi_2) \end{aligned} \quad (34)$$

or it can be written

$$\begin{aligned} \kappa_{12}(T, p + \pi_2, n_1, n_2)V_{12}(T, p + \pi_2, n_1, n_2) \\ = n_1\bar{V}_1^0(T, p - \pi_1 + \pi_2)\kappa_1^0(T, p - \pi_1 + \pi_2) + n_2\bar{V}_2^0(T, p)\kappa_2^0(T, p). \end{aligned} \quad (35)$$

Again, the volumes and compressibilities of the solution and of the pure constituents can be measured at the pressures indicated in Eqs. (34) and (35). It should be noted that  $n_1$  and  $n_2$  refer to the molar amounts of 1 and 2 at  $T$  and  $p + \pi_1$  in Eq. (34) and to the molar amounts at  $T$  and  $p + \pi_2$  in Eq. (35). In many solutions,  $n_1$  will differ from  $n_1^0$  and  $n_2$  will differ from  $n_2^0$  at the same  $T$  and  $p$ . Furthermore,  $n_1$  and  $n_2$  may vary with  $T$  and  $p$ . Nevertheless, Eqs. (33), (34) and (35) are exact relationships and they provide an opportunity to investigate  $n_1$  and  $n_2$  in various solutions at various temperatures and pressures.

According to Eq. (33) an increase in the pressure applied to the solution from  $p$  to  $p + \Delta p$ , an increase from  $p - \pi_1$  to  $p + \Delta p - \pi_1$  applied to the pure solvent and an increase from  $p - \pi_2$  to  $p + \Delta p - \pi_2$  applied to the pure solute decreases their volumes at constant  $T$

$$\Delta V(T, p, n_1, n_2) = n_1 \Delta \bar{V}_1^0(T, p - \pi_1) + n_2 \Delta \bar{V}_2^0(T, p - \pi_2). \quad (36)$$

Similarly, increasing the pressures in Eqs. (34) and (35) by  $\Delta p$  decrease the volumes of solution, pure solvent and pure solute the amounts

$$\Delta V(T, p + \pi_1, n_1, n_2) = n_1 \Delta \bar{V}_1^0(T, p) + n_2 \Delta \bar{V}_2^0(T, p + \pi_1 - \pi_2) \quad (37)$$

and

$$\Delta V(T, p + \pi_2, n_1, n_2) = n_1 \Delta \bar{V}_1^0(T, p + \pi_2 - \pi_1) + n_2 \bar{V}_2^0(T, p). \quad (38)$$

Thus, the variation for  $n_1$  and  $n_2$  as a function of  $T$  and pressure can be experimentally determined for a solution constituted by mixing  $n_1^0$  moles of pure solvent and  $n_2^0$  moles of pure solute.

## References

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